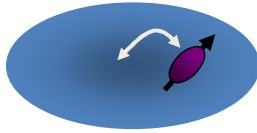


Quadrupolar coupling

theory

Quadrupolar couplings



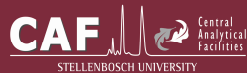
= electric interaction of spin $>1/2$ nuclei with the surrounding electric fields

Very often as strong as Zeeman interaction!

interaction between a quadrupole moment (Q) with an electric field gradient (EFG)

→ Electric field gradient (EFG): ???

→ Electric quadrupole moment: ???



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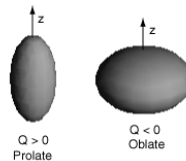
Quadrupolar coupling

Theory: the electric quadrupole moment Q

- Feature of all nuclei with spin $>1/2$:

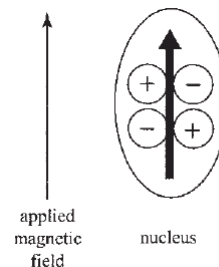
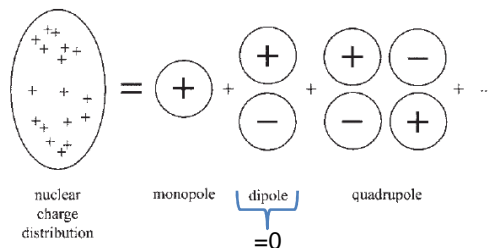
non-spherical distribution of the nuclear charge

(75% of all NMR active nuclei)



The quadrupole moment is fixed within the nucleus:
→ alignment along B_0 field through the spin

- The non-spherical charge distribution can be described mathematically by a multipole expansion:



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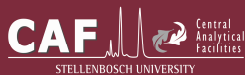
Quadrupolar coupling

Theory: the electric quadrupole moment Q

Table 9.1 Quadrupole moments of selected nuclear isotopes with $I > 1/2$. A complete listing of nuclear quadrupole moments may be found on the website www.webelements.com.

Isotope	Ground-state spin	Natural abundance/%	Electric quadrupole moment/ 10^{-28}m^2
^2H	1	0.012	0.2860
^6Li	1	7.59	-0.0808
^7Li	3/2	92.41	-4.01
^{11}B	3/2	80.1	4.059
^{14}N	1	99.6	2.044
^{17}O	5/2	0.038	-2.558
^{23}Na	3/2	100	10.4
^{27}Al	5/2	100	14.66
^{45}Sc	7/2	100	-22
^{51}V	7/2	99.8	-5.2
^{55}Mn	5/2	100	0.33
^{59}Co	7/2	100	0.42
^{63}Cu	3/2	69.2	-0.22
^{65}Cu	3/2	30.8	-0.204
^{87}Rb	3/2	27.8	0.132
^{93}Nb	9/2	100	-0.32

!!! The strength of quadrupolar coupling depends on Q AND the EFG!



M. Levitt, spin dynamics, 2nd edition, Wiley & sons, (2008)

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Quadrupolar coupling

Theory: the electric field gradient (EFG)

- arises if **electron density is non-spherical distributed**
= feature of the electronic environment (depends on compound, morphology...)

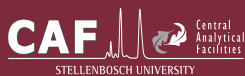
- Mathematically described by an tensor (like the chemical shift tensor):

$$\mathbf{V} = \begin{pmatrix} V_{xx} & V_{xy} & V_{xz} \\ V_{yx} & V_{yy} & V_{yz} \\ V_{zx} & V_{zy} & V_{zz} \end{pmatrix}$$

$$\mathbf{V}^{PAF} = \begin{pmatrix} V_{XX}^{PAF} & 0 & 0 \\ 0 & V_{YY}^{PAF} & 0 \\ 0 & 0 & V_{ZZ}^{PAF} \end{pmatrix}$$

The trace of the EFG tensor is zero:

$$V_{XX} + V_{YY} + V_{ZZ} \\ = V_{XX}^{PAF} + V_{YY}^{PAF} + V_{ZZ}^{PAF} = 0$$



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Quadrupolar coupling

Theory: parameters of the EFG tensor

Analog calculation like for the chemical shift tensor:

- isotropic component: $V^{iso} = \frac{1}{3}(V_{xx} + V_{yy} + V_{zz}) = 0$ (because the trace of the EFG tensor is zero)

→ no isotropic component of the EFG tensor

- anisotropy: $V^{aniso} = V_{zz}^{PAF} - V^{iso} = V_{zz}^{PAF} - 0 = V_{zz}^{PAF}$ with: $V_{zz}^{PAF} = eq$

→ anisotropy of the EFG tensor is the largest component of the vector

- asymmetry: $\eta_Q = \frac{V_{yy}^{PAF} - V_{xx}^{PAF}}{V_{zz}^{PAF}}$ $0 \leq \eta_Q \leq 1$

Quadrupolar coupling

Theory: the quadrupolar Hamiltonian

In the PAF: $\hat{H}_Q = \underbrace{\frac{eq \cdot eQ}{h}}_{\text{Quadrupolar coupling constant [rad/s]}} \cdot \frac{1}{4I(2I-1)} \left[3\hat{I}_{z^{PAF}}^2 - \hat{I}^2 + \frac{1}{2}(\hat{I}_{x^{PAF}}^2 - \hat{I}_{y^{PAF}}^2) \right]$

Quadrupolar coupling constant [rad/s]:

Quadrupolar coupling constant [Hz]: $C_Q = \frac{eq \cdot eQ}{h}$

→ Measure of the strength for the quadrupolar interaction

Quadrupolar coupling

Theory: impact on the NMR energy level

So far:

- all interactions were considered as perturbations of the Zeeman interaction:

$$\hat{H} = \hat{H}_Z + \hat{H}_X$$

this perturbation theory approach is valid, if the interaction X is small compared to the Zeeman interaction

→ for quadrupolar coupling very often not the case!

Perturbation theory must be expanded:

$$\hat{H} = \hat{H}_Z + \hat{H}_Q^{(1)} + \hat{H}_Q^{(2)} + \dots$$

$E_Q^{(1)}$

$E_Q^{(2)}$

"first order/ second order energy correction term of the quadrupolar coupling interaction"

Quadrupolar coupling

Theory: impact on the NMR energy level

Energy correction terms of

- First order:

Can be averaged out by MAS!

$$E_{Q_m}^{(1)} = \frac{eq \cdot eQ}{4I(2I-1)} \left(3m^2 - I(I+1) \right) \frac{1}{2} \left[(3 \cos^2 \theta - 1) + \eta_Q \cos 2\phi \sin^2 \theta \right]$$

= $C_Q h$

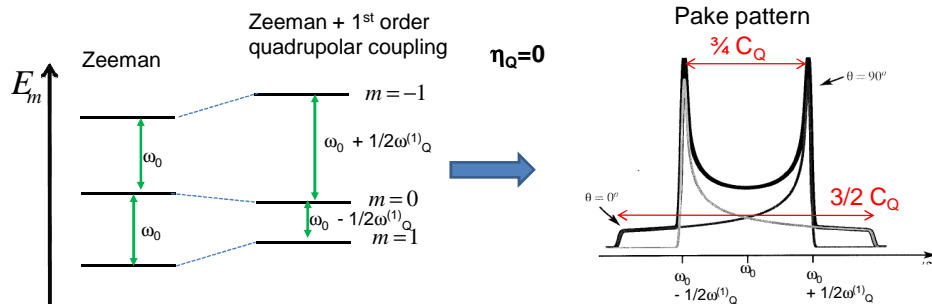
Same energy shifts

for +m and -m!

→ Central transition ($m=1/2 \leftrightarrow m=-1/2$) is not affected!

Quadrupolar coupling

Spin-1 nuclei: energy diagrams and lineshapes



- Only natural $I=1$ spins: ^2H , ^6Li and ^{14}N

- ^6Li : - nuclei with smallest Q

- usually Pake pattern can't be observed due to distribution of C_Q
 → Gaussian lineshapes

- ^{14}N : - small Q , but often strong EFGs → quadrupolar coupling in order of several MHz
 → hard to impossible to measure

Quadrupolar coupling

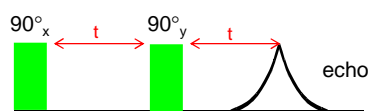
Spin-1 nuclei: ^2H NMR

→ Low Q , C_Q s typically about 130 kHz

→ correspond to an FID decay within 5 μs (dead time about 4 μs)

→ ^2H NMR spectra can be measured using the Solid Echo pulse sequence

The Solid Echo (quadrupole echo) pulse sequence:



- designed for spin-1 nuclei
 - can refocus 1st order quadrupolar coupling and homonuclear dipolar coupling

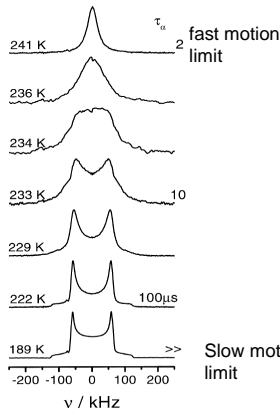
- Variation for spin-3/2 nuclei:
 $90^\circ_x - t - 64^\circ_y - t - \text{aq}$

Quadrupolar coupling

²H NMR: study of dynamics in solids

²H NMR

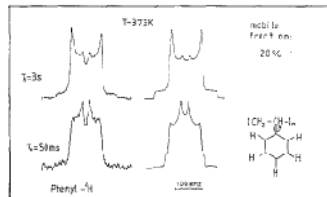
glycerol-d₅



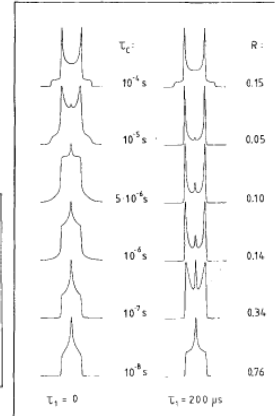
motional narrowing:

→ advantage of ²H NMR:

Differentiation of motion types !



²H spectra of PS-d₅ showing the influence of 180° phenyl flips



calculated ²H spectra for two site jumps



Spies, H.W.; Colloid & Polymer Sci., 256, 193-209, (1983)

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Quadrupolar coupling

Spin-3/2 nuclei: energy diagrams and lineshapes

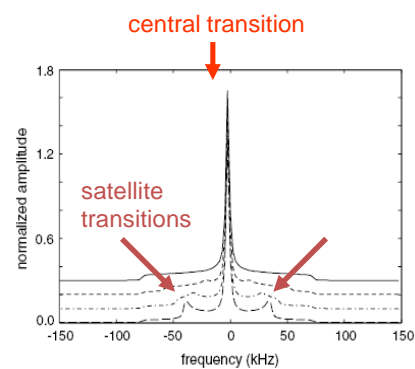
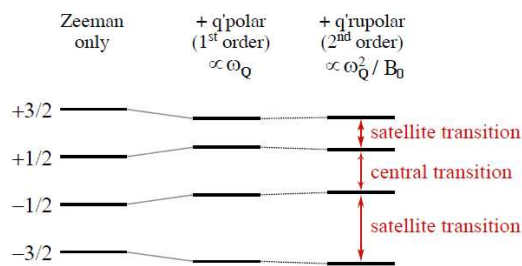


Fig. 6. Simulations of powder pattern spectra for spin $I=3/2$ for a coupling constant $e^2qQ/h = 100$ kHz and various asymmetry parameters. To simulate the dipolar broadening of the pattern the spectra were convoluted with a 100 Hz wide Lorentzian line-shape. The spectra were shifted vertically for clarity. From top to bottom η was 1, 0.5, 0.2, and 0

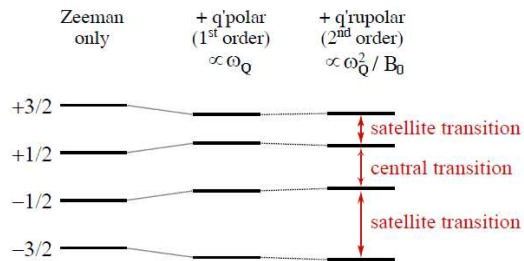


R. Böhmer, K.R. Jeffrey, M.Vogel;
J. Prog. Nucl. Magn. Reson. Spectrosc., 50, 87-174, (2007)

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Quadrupolar coupling

Spin-3/2 nuclei: hard and soft pulses



Homogenous excitation:

- Central as well satellite transitions are excited
- Requires short pulses with high amplitudes ("hard pulse")
- $90^\circ(\text{solid}) = 90^\circ(\text{solution})$

Selective excitation:

- Only central transition is excited
- requires long pulses with a small amplitude ("soft pulse")
- $90^\circ(\text{solid}) = 1/2 \times 90^\circ(\text{solution})$

**In both cases:
sinus-nutation curve!**

Quadrupolar coupling

Spin-3/2 nuclei: hard and soft pulses

Should I use hard or soft pulses?

→ Depends on the purpose of the measurement!

Example:

^{23}Na Spin Echo Decay spectroscopy requires selective excitation!
(to study average ^{23}Na - ^{23}Na distances)

Quantitative one pulse MAS experiments require a homogenous excitation!

Quadrupolar coupling

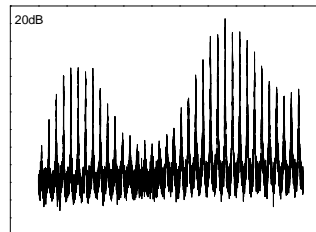
Spin-3/2 nuclei: hard and soft pulses

Non-symmetrical nutation curve without negative intensities!

???

→ Not homogenous,
not selective excitation either!

nutation curve:



-Central transition and satellite transitions
are partially excited.

Try to avoid this scenario!

- Increase pulse power level → homogenous exc.
- decrease pulse power level → selective exc.

In real world: sometimes not achievable!

→ Use pulses with the length: $90^\circ/(I+1/2)$

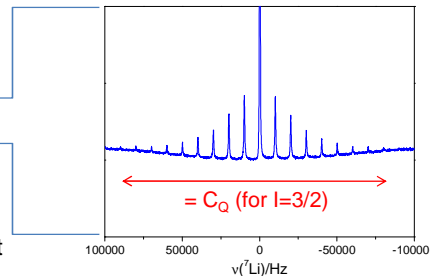
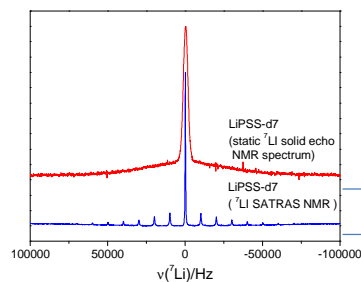
Quadrupolar coupling

MAS experiments: weak quadrupolar coupling

Satellite Transition spectroscopy (SATRAS)

= MAS experiment using hard pulses and big spectral width (and many scans)

→ Aim: complete excitation of all satellite transitions

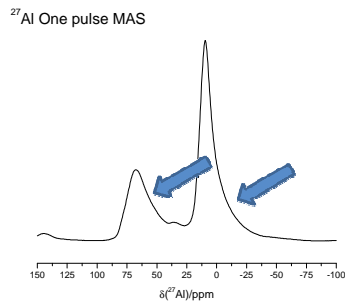


→ C_Q can be estimated from most
outer spinning sidebands

Quadrupolar coupling

MAS experiments: strong quadrupolar coupling

- First and second order quadrupolar coupling perturbations
- Interpretation of one pulse MAS NMR spectra complicated



Reason: affected by 2nd order quadrupolar interaction

→ Isotropic chemical shift can't be determined from peak position

Typical: asymmetric lineshape (due to distribution of quadrupolar coupling constants)

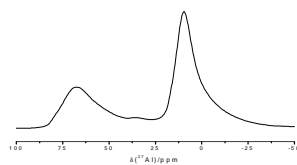
→ could cover different positions!

Quadrupolar coupling

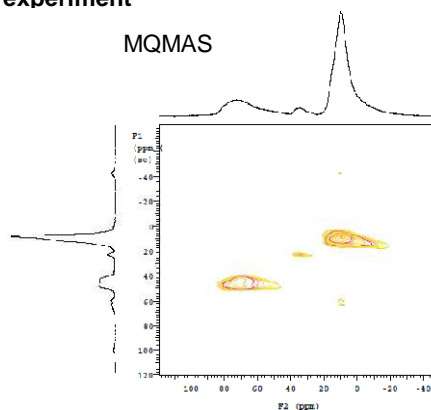
MAS experiments: strong quadrupolar coupling

The multiple-quantum MAS (MQMAS) experiment

one pulse MAS



MQMAS



MQMAS :
2D experiment to obtain
high-resolution spectra of half-integer
quadrupolar nuclei

Some wise, closing words... 😊

Before you start any Solid State NMR experiments, please ... :

-Be aware that solution state and solid state NMR are very different techniques (different strengths and weaknesses)

-Think about:

What is your research question?

→ What experiments could be useful to answer this question?

Accept that you won't find a recipe to interpret Solid State NMR results!

→ Interpretation requires understanding of

- solid state NMR (interactions, their impact on the spectra, methods)

- your material (what nuclei, nuclei environment, morphology,...)

→ Practice makes perfect! 😊

I hope this workshop was a first step in understanding Solid State NMR!

Thank you for your attention!

Quadrupolar coupling

Theory: impact on the NMR energy level

Energy correction terms of

- Second order:

- no orientation dependence
→ isotropic component

$$E_{Q_m}^{(2)} = - \left(\frac{eq \cdot eQ}{4I(2I-1)} \right)^2 \frac{m}{\omega_0} \times \left\{ -\frac{1}{5} (I(I+1) - 3m^2) (3 + \eta_c^2) \right.$$

$$+ \frac{1}{28} (8I(I+1) - 12m^2 - 3) (\eta_c^2 - 3) (3 \cos^2 \theta - 1) + 6\eta_c \sin^2 \theta \cos 2\phi \left. \right]$$

$$+ \frac{1}{8} (18I(I+1) - 34m^2 - 5) \left[\frac{1}{140} (18 + \eta_c^2) (35 \cos^4 \theta - 30 \cos^2 \theta + 3) \right.$$

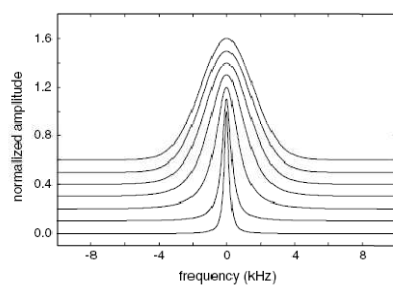
$$\left. + \frac{3}{7} \eta_c \sin^2 \theta (7 \cos^2 \theta - 1) \cos 2\phi + \frac{1}{4} \eta_c^2 \sin^4 \theta \cos 4\phi \right\}$$

-Increases with $(C_Q/h)^2$

- Energy shifts in opposite directions for +m and -m!

-Proportional to $1/B_0$
→ 2nd order quadrupolar coupling is reduced at strong magnetic fields!!

-Complicated orientation dependence
→ can't be averaged out by MAS!



slow motion limit:
Gaussian line

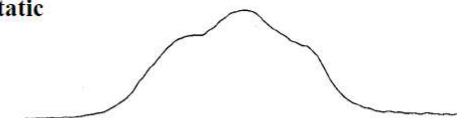
motional narrowing:
superposition of a Gaussian and
a Lorentzian line

fast motion limit:
Lorentzian line

MQMAS in action

^{23}Na spectra (105.8 MHz)
of sodium citrate
S. C. Wimperis

Static



MAS

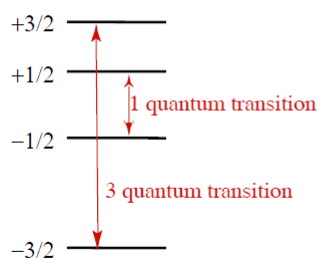


MQMAS



N.B. Horizontal scale varies

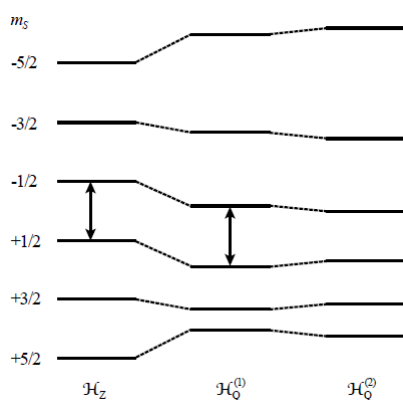
$I=3/2$



- 1Q and 3Q transition frequencies are unaffected by 1st order quadrupole interaction
- Contributions of isotropic and 2nd order components are non-zero but *different*
- Multiple Quantum Magic-Angle Spinning (MQMAS) correlates 1Q and 3Q frequencies to separate isotropic and 2nd order effects
- 2D experiment that suppresses quadrupole broadening with standard MAS hardware!

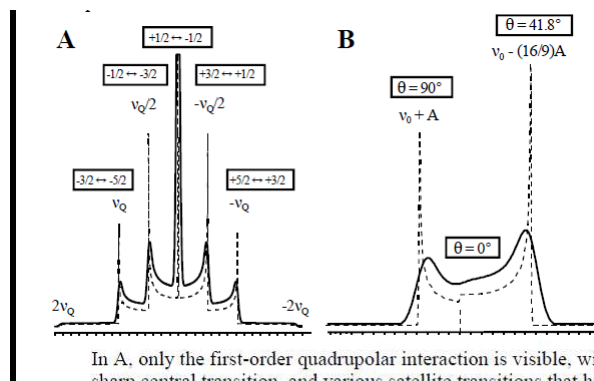
Solid state NMR has been applied to:

organic complexes	inorganic complexes
zeolites	mesoporous solids
microporous solids	aluminosilicates/phosphates
minerals	biological molecules
glasses	cements
food products	wood
ceramics	bones
semiconductors	metals and alloys
archaeological specimens	polymers
resins	surfaces



Solid state NMR has been applied to:

organic complexes	inorganic complex
zeolites	mesoporous solids
microporous solids	aluminosilicates/p
minerals	biological molecu
glasses	cements
food products	wood
ceramics	bones
semiconductors	metals and alloys
archaeological specimens	polymers
resins	surfaces



In A, only the first-order quadrupolar interaction is visible, with a sharp central transition, and various satellite transitions that have shapes resembling axial CSA patterns.

In B, the value of C_Q is much larger. The satellite transitions broaden and disappear and only the central transition spectrum is left (which is unaffected by first-order interactions). It still has a strange shape due to the orientation dependence of the second-order quadrupolar frequency.



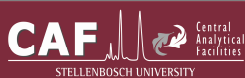
repetitions (constant experiment time)	recycle / T_1	relative S/N
8	5	1.00
12	3.33	1.19
32	1.2	1.41
80	0.5	1.25
400	0.1	0.68

$$\text{tip-angle} = \cos^{-1}(e^{-t/T_1})$$

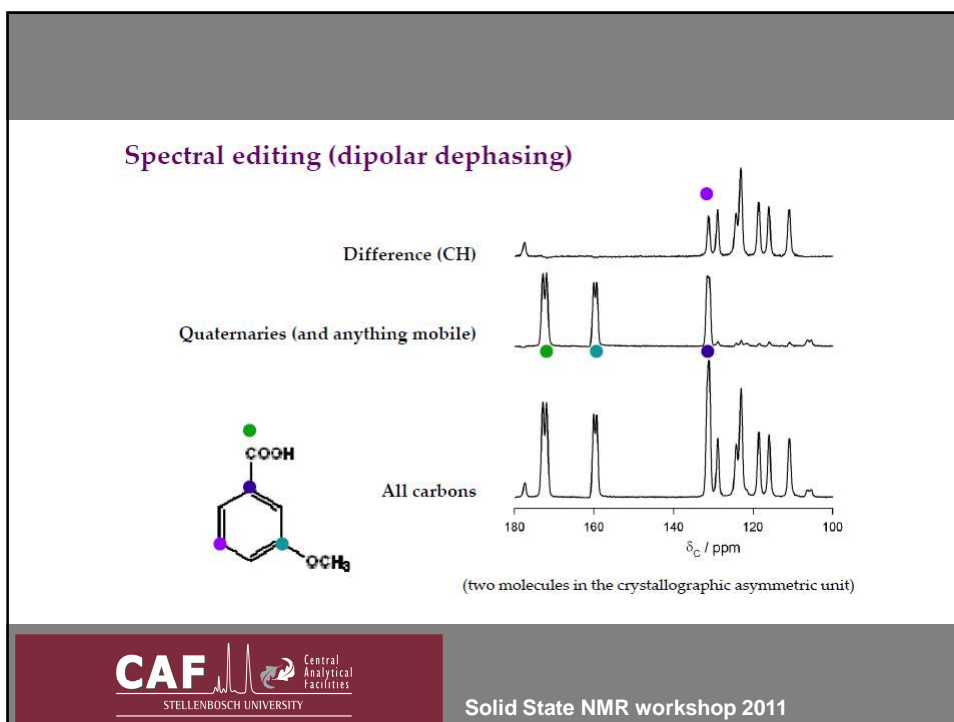
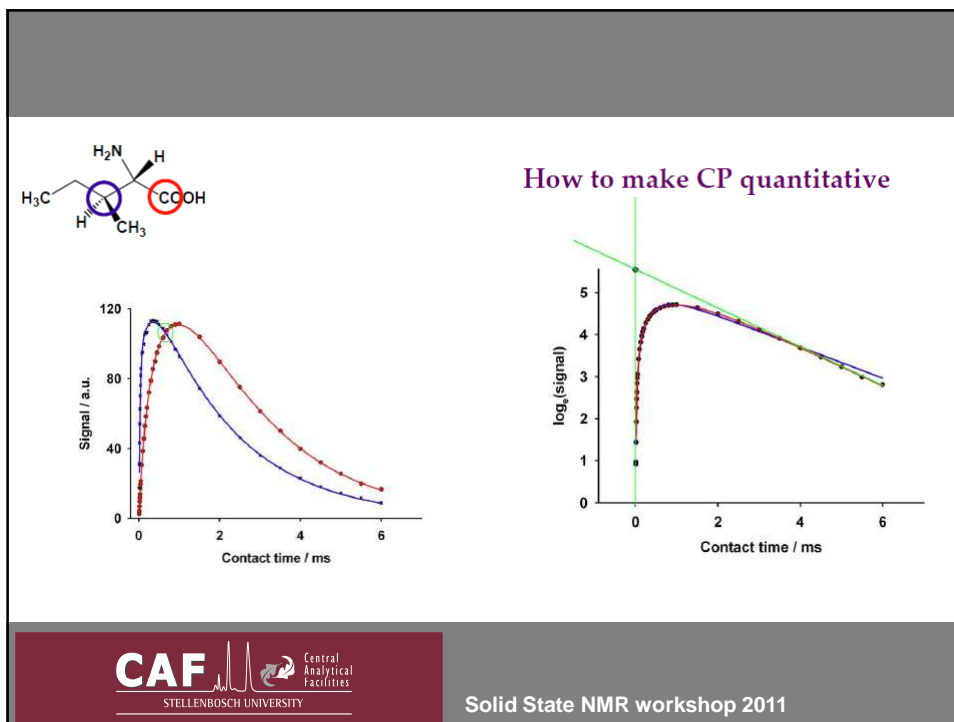
“Ernst” angle

$$\text{tip-angle} < 90/(I+1/2)$$

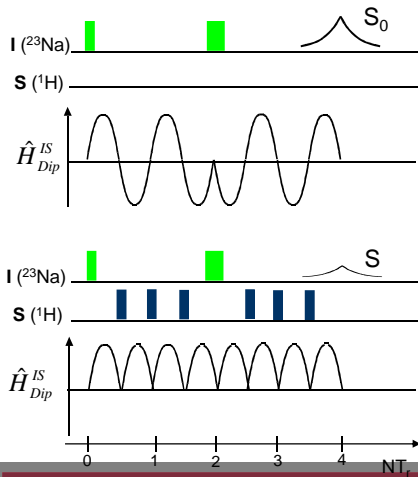
quadrupoles



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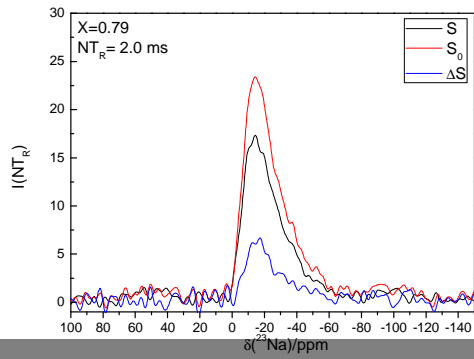


The REDOR pulse sequence:

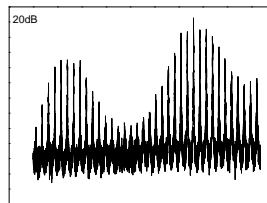


REDOR= Rotational Echo Double
Resonance

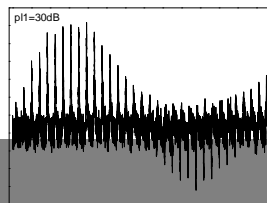
$$\hat{H}_{Dip}^{IS} \propto \frac{\hat{I}_Z \hat{S}_Z}{r^3} (3 \cos^2 \theta - 1)$$



nutations curves:



no selective excitation of the
central transition !



53

In this [shear mapping](#) the red arrow changes direction but the blue arrow does not. Therefore the blue arrow is an eigenvector, with eigenvalue 1 as its length is unchanged.

